

## CHARACTERIZATION OF LASER ANNEALED SE-IMPLANTED GaAs

A. H. ORABY\*

*SUMMARY: A Q-switched ruby or Nd: Glass laser has been used to anneal Se-implanted GaAs. Differential Hall-effect and sheet-resistivity measurements have been carried out to investigate the electrical properties of the implanted layers as a function of laser energy density and laser wavelength. Electron concentrations greater than  $1 \times 10^{19}/\text{cm}^3$  were obtained but mobility values were low. Significant indiffusion of Se atoms was observed following laser irradiation. The depth of electrically active layer greatly depends on the laser energy density. A comparison of the measured annealed depth as a function of laser energy density with the theoretically calculated annealed depth is presented.*

*Key Words: Laser, Laser annealing.*

### INTRODUCTION

Since the pioneering work of laser annealing (1-3) of ion implantation damage in silicon and gallium arsenide, there have been a rapidly growing effort devoted to use laser annealing for the removal of implantation damage and the activation of implanted species in GaAs and other semiconductors (4-8). Pulse laser annealing is a very attractive processing technique for GaAs and other semiconductors as a means of capless annealing. However, there are problems in the pulsed laser annealing of GaAs, such as the lack of activation in low dose implanted GaAs and low mobilities in high dose implanted samples (9-11).

The efficiency of the laser annealing depends mainly on the absorption coefficient, which in turn depends on the wavelength and the material structure. Two different models have been suggested to explain pulsed laser annealing. The first model deals with a thermal transport model to describe the melting and resolidification of semiconductors (12-17). The other model supposes that pulsed laser annealing is primarily a nonthermal effect, resulting from formation of a dense electron-hole plasma (18).

In the present work, the dependence of pulsed laser annealing effects on the laser wavelength and energy density are investigated. Comparison of the electrical properties of the implanted GaAs between laser annealed and thermally annealed samples is represented. Comparison of experimental and theoretical melt depth produced by pulsed laser irradiation as a function of laser energy density are discussed.

### EXPERIMENTAL PROCEDURE

Semi-insulating GaAs wafers of (100) orientation were implanted at room temperature with 60 keV Se ions at a dose of  $1 \times 10^{15}/\text{cm}^2$  in a non-channeling direction. Laser annealing was performed in air at room temperature with a Q-switched ruby laser ( $\lambda=0.694 \mu\text{m}$ ) or Nd: Glass laser ( $\lambda=1.06 \mu\text{m}$ ) with a single pulse of 20 nsec duration (FWHM) without using an encapsulant. The laser energy densities employed were 0.2-0.7 J/cm<sup>2</sup>. A lapped quartz diffuser rod was used between the laser and the sample to obtain a homogeneous intensity distribution at the sample surface. For comparison purpose, some samples were thermally annealed in a flowing H<sub>2</sub> at 850°C for 20 min after encapsulation with approximately 1500-Å Si<sub>3</sub>N<sub>4</sub> film.

---

\*From Department of Physics, Faculty of Science, Mansoura University, Mansoura, Egypt.

Table 1: Sheet electron concentration ( $n_s$ ), sheet mobility ( $\mu_s$ ) and the percentage of sheet electron concentration to the implantation dose ( $n_s/D$ ) for 60 keV Se implanted GaAs at a dose of  $1 \times 10^{15}/\text{cm}^2$  and annealed at three different conditions.

Property	Thermal annealing 850°C/20 min	Laser annealing (0.6 J/cm <sup>2</sup> )	
		Ruby laser	Nd: Glass laser
$n_s$	$5 \times 10^{13}$	$1.4 \times 10^{14}$	$1.1 \times 10^{14}$
$\mu_s$	1900	580	560
$n_s/D\%$	5	14	11

Electron concentration and mobility profiles were obtained by differential Hall effect and sheet resistivity measurements combined with successive anodization and layer removal technique (19). Ohmic contacts were formed by deposited Au-Ge (12% Ge weight percent) dots. In the thermally - annealed samples, the contacts were alloyed for 2 min in flowing H<sub>2</sub>. In the laser annealed samples, the contacts were formed by irradiating the back surface of the samples with a Q-switched Nd: Glass laser (20) at an energy density of 0.45 J/cm<sup>2</sup>, since heat treatment above 200°C after laser annealing reduces the carrier concentration (21).

#### RESULT AND DISCUSSION

A comparative summary of the results of Hall-effect measurements made on Se-implanted GaAs, annealed at three different conditions is given in Table 1, where the values of sheet electron concentration  $n_s$ , electron mobility  $\mu_s$  and the percentage of sheet electron concentration to the implantation dose ( $n_s/D$ ) are tabulated. The measurements were made after etching van der Pauw-type mesas provided with Au-Ge ohmic con-

tacts (22). The data in Table 1 show comparable values of  $n_s$ ,  $\mu_s$  and the electrical activity ( $n_s/D$ ) for both the samples annealed by the laser pulses of Ruby or Nd: Glass. The laser annealed samples showed higher values of  $n_s$  and lower values of  $\mu_s$  than those obtained in thermally annealed sample. The relatively low values of  $\mu_s$  for laser annealed samples is an indication of the presence of defects in these samples, even taking into account the high concentration of activated Se ions (8,9,11). The laser annealed samples showed a better electrical activity than the thermally annealed samples.

Table 2 shows the sheet electron concentration  $n_s$ , the electron mobility and the percentage of electrical activity ( $n_s/D$ ) of 60 keV Se-implanted GaAs as a function of laser energy density and laser wavelengths. The data shows that the doping efficiency or the electrical activity increases with increases in laser energy density. As shown, pulsed ruby laser annealing gives better electrical activation values than pulsed Nd: Glass laser in Se-implanted GaAs. The electrical properties greatly depends on laser energy density. For

Table 2: Electrical properties of 60 keV Se-implanted GaAs at a dose of  $1 \times 10^{15}/\text{cm}^2$  and pulsed laser annealed by ruby or Nd: Glass laser and different energy densities.

Laser Energy (J/cm <sup>2</sup> )	Ruby laser			Nd: Glass laser		
	$n_s$ (cm <sup>-2</sup> )	$\mu_s$ (cm <sup>2</sup> /V.s)	$n_s/D\%$	$n_s$ (cm <sup>-2</sup> )	$\mu_s$ (cm <sup>2</sup> /Vs)	$n_s/D\%$
0.3	$1.2 \times 10^{13}$	390	1.2	$9.2 \times 10^{12}$	400	0.92
0.4	$5.5 \times 10^{13}$	420	5.5	$3.2 \times 10^{13}$	430	3.2
0.5	$9.5 \times 10^{13}$	480	9.5	$6.5 \times 10^{13}$	510	6.5
0.6	$1.4 \times 10^{14}$	580	14	$1.1 \times 10^{14}$	560	11

samples annealed with laser energy density less than 0.2 J/cm<sup>2</sup>, no activity of Se-implanted GaAs was obtained. The threshold energy density for electrical activation in both laser types was around 0.2 J/cm<sup>2</sup>, but the electrical properties was poor at this energy. This indicates that this energy density was not enough to melt all of the entire amorphous layer, thus leading to polycrystalline regrowth on an underlying heavily damaged material (23). A visible surface damage was observed at an energy density greater than 0.65 J/cm<sup>2</sup>. The results of Table 2 shows that Q-switched ruby laser

is more effective than Q-switched Nd: Glass laser for annealing Se-implanted GaAs to obtain high electrical activation.

Figure 1 shows depth profiles of the electron concentration and mobility-measured on samples implanted with 60 keV Se ions at a dose of 1x10<sup>15</sup>/cm<sup>2</sup> and laser annealed with Nd: Glass laser at three different energy densities. The Se atomic profile of the as-implanted sample predicted by LSS theory (24) is also plotted for comparison. The profiles show the diffusion (25) of Se atoms at laser energy densities of 0.5 and

Figure 1: Depth profiles of the electron concentration and mobility measured on samples implanted with 60 keV Se ions at a dose of 1 x 10<sup>15</sup>/cm<sup>2</sup> and annealed at three different energy densities with Q-switched Nd: Glass laser.

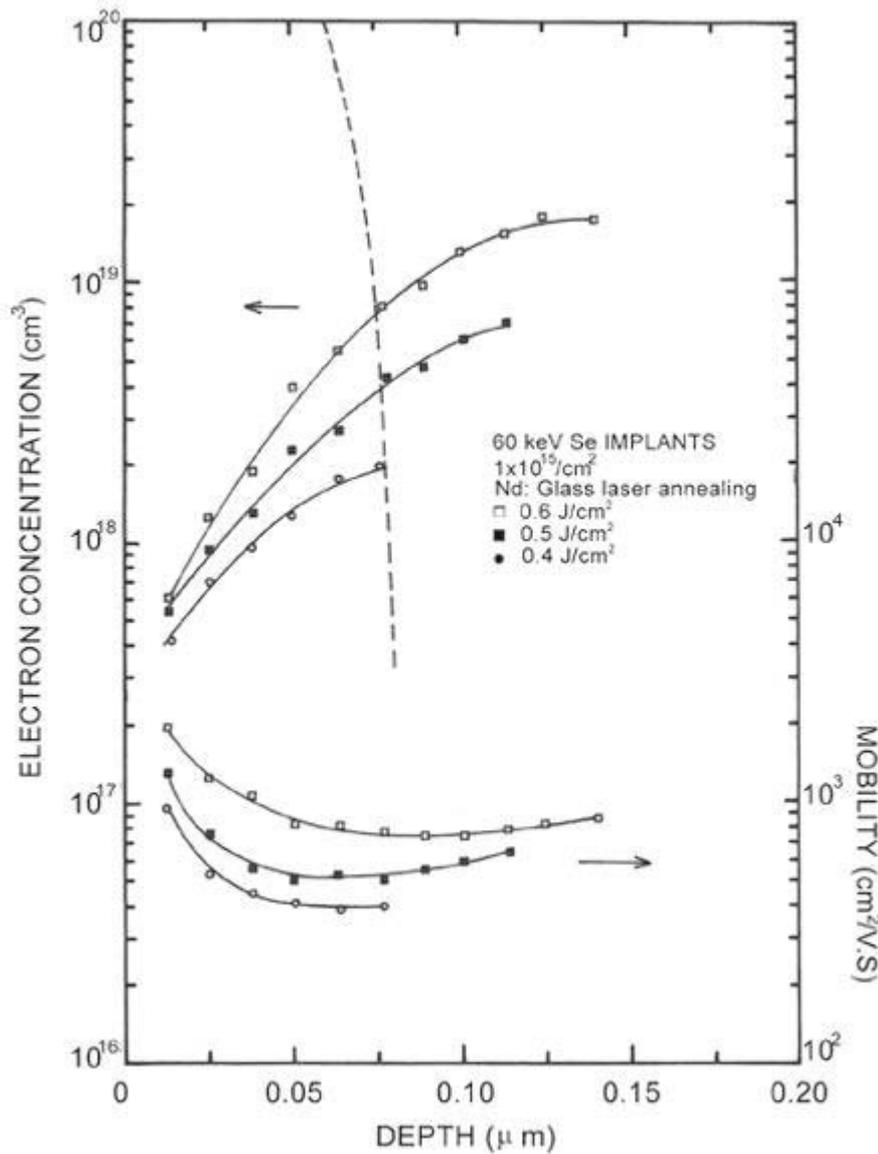
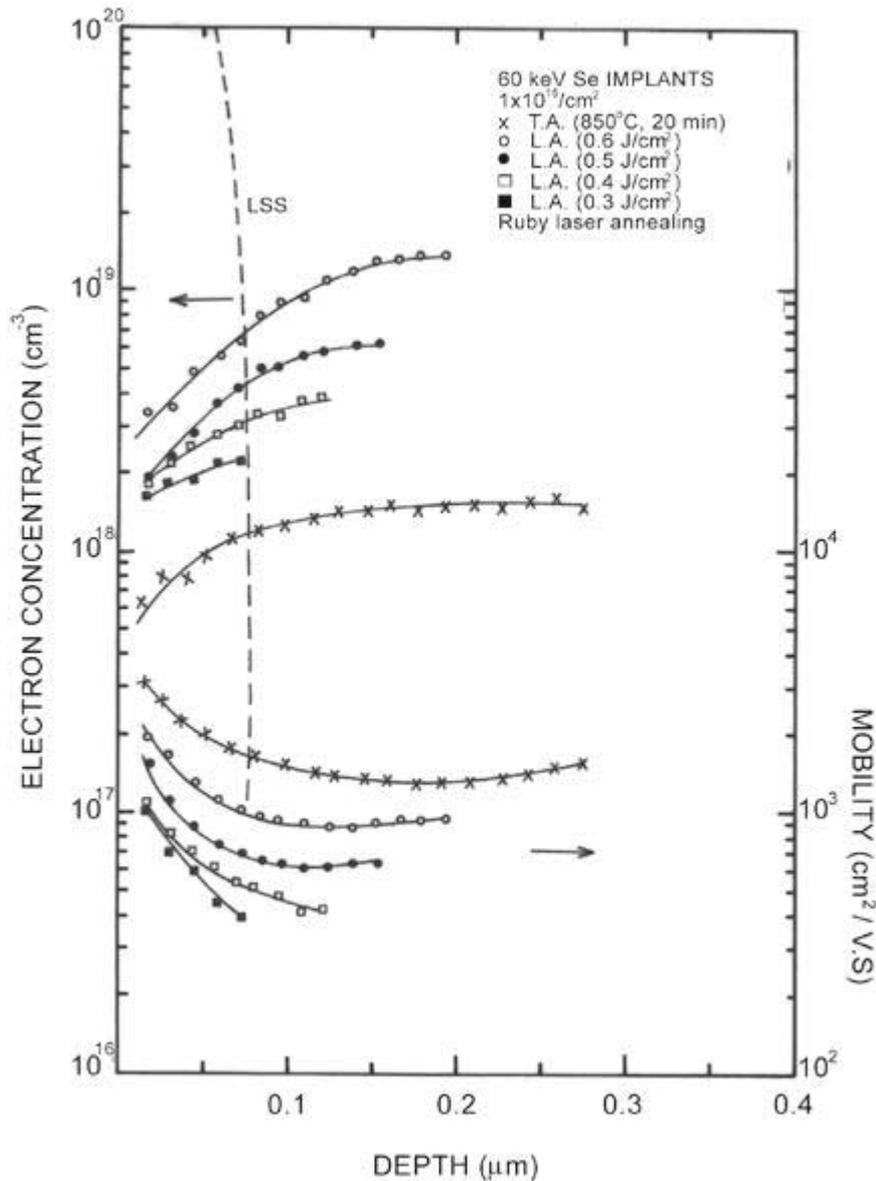


Figure 2: Depth profiles of the electron concentration and mobility measured on samples implanted with 60 keV Se ions at a dose of  $1 \times 10^{15}/\text{cm}^2$  and annealed Q-switched ruby laser at four different energy densities or thermally annealed.



0.6 J/cm<sup>2</sup> compared with that predicted by LSS theory. For the sample annealed at an energy density of 0.4 J/cm<sup>2</sup>, the highest electron concentration obtained was  $1.9 \times 10^{18}/\text{cm}^3$  with mobility of  $\sim 400 \text{ cm}^2/\text{V.s}$ , whereas the electrically active layer extended to a depth of 800-Å from the surface. For sample annealed at an energy density of 0.6 J/cm<sup>2</sup>, the highest electron concentration obtained was  $1.8 \times 10^{19}/\text{cm}^3$  with mobility of  $\sim 550 \text{ cm}^2/\text{V.s}$ , whereas the electrically active layer extended to a depth of 1400-Å from the surface. The

profiles show that there is a less conductive layer near the surface which may be attributed to laser induced defects and in-diffusion of Se atoms. The electron concentration profiles shows a great dependence on the laser energy density, whereas the highest electron concentration was increased by approximately one order with increasing laser annealing energy density from 0.4 to 0.6 J/cm<sup>2</sup>.

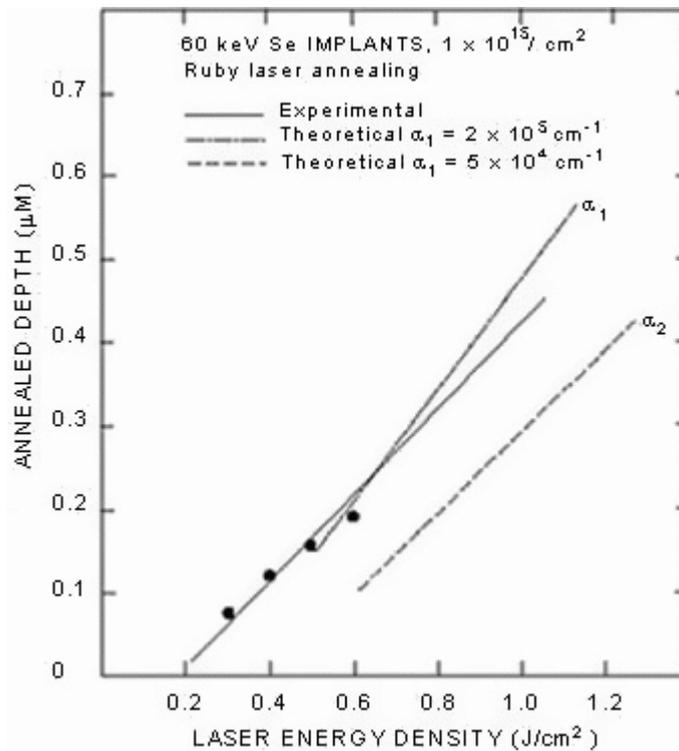
Figure 2 shows electron concentration and electron mobility profiles measured on samples implanted at a

dose of  $1 \times 10^{15}/\text{cm}^2$  Se ions and annealed at four different laser energy densities of a pulsed ruby laser compared with the conventional thermal annealing. The LSS profile is plotted for comparison. For a sample, thermally annealed at  $850^\circ\text{C}$  for 20 min, the highest electron concentration is  $1.5 \times 10^{18}/\text{cm}^3$  with mobility  $1400\text{cm}^2/\text{V.s}$ , and the measured electron concentration profiles are deeper than the prediction of Se penetration based on LSS range theory. This penetration seems to be due to significant diffusion (25,26) of Se-implanted ions. For samples subjected to pulsed ruby laser irradiation at different energy densities, both electron concentration and electron mobility profiles show a great dependence on the laser energy density. For samples annealed with  $0.6 \text{ J}/\text{cm}^2$  a high electron concentration is obtained of  $1.4 \times 10^{19}/\text{cm}^3$  which exceeds that of the thermally annealed samples by about one order. After laser annealing, the electron concentration profile shows a peak shift from the surface deeper than the depth predicted by LSS theory. This behaviour of Se in GaAs has been reported in pulsed laser annealing (27-28). This shift may be due to the effect of laser

induced defects near the surface region which partially decreases the electron concentration near the surface and is also due to the redistribution of implanted Se in GaAs.

Figure 3 shows the experimental and theoretical values of the annealed depth as a function of ruby laser energy density for 60 keV SE-implanted GaAs at a dose of  $1 \times 10^{15}/\text{cm}^2$ . Each point in the experimental data represents the thickness from the surface of the last possible measurements before reaching high resistance material with respect to the laser energy density employed. The position of these points from the surface are considered to be equal the depth of the melted layer that depends on the laser energy density. These points lie in a straight (solid line) intercept with the laser energy axis around the value of  $0.18 \text{ J}/\text{cm}^2$ . This value of energy density represents the threshold energy density for melting the surface of the implanted layer. This value is in agreement with Rutherford backscattering measurements (8,29), that shows no change in the crystal quality from that of the as-implanted sample for laser energy density less than 0.2

Figure 3: Experimental and theoretical values of the annealed depth as a function of laser energy density for 60 keV Se-implanted GaAs at a dose of  $1 \times 10^{15}/\text{cm}^2$ .



J/cm<sup>2</sup>. The theoretical (dashed lines) melted depth as a function of laser energy density was obtained from the reported data of Wang *et al.* (14). The calculation was based on the thermal melting model of pulsed ruby laser annealing of GaAs. The variation of the annealed depth (melted depth) with the laser energy density and the absorption coefficient is plotted for the purpose of comparison with the experimental data which shows some agreements with the theoretical calculation based on the value of the absorption coefficient  $\alpha_1=2 \times 10^5 \text{ cm}^{-1}$ .

### CONCLUSIONS

The effects of pulsed laser annealing on the electrical properties of Se-implanted GaAs as a function of laser energy density and laser wavelength were investigated by means of differential Hall effect measurements. Laser annealing of the sample implanted at a dose of  $10^{15}/\text{cm}^2$  yielded an electron concentration higher than that obtained by thermal annealing and a lower electron mobility than that expected in uncompensated bulk GaAs. Electrical properties of Se-implanted GaAs greatly depend on laser energy density. There is a threshold laser energy density for the recovery of implantation damage and electrical activation of the implanted ions. The laser annealing energy density should be enough to melt the entire implanted layer.

### ACKNOWLEDGEMENTS

The author would like to thank Prof. K. Gamo, and Mr. K. Kawasaki of Faculty of Engineering Science, Osaka University, Osaka, Japan, for their assistance during ion implantation and laser annealing.

### REFERENCES

1. Shtyrkov EI, IB Khaibullin, MM Zaripov, MF Galyatudinov and RM Bayazitov : *Sov Phys Semicond* 9, p 1309, 1975.
2. Kachurin GA, NB Pridachin and LS Smirnov : *Sov Phys Semicond* 9, 946, 1975.
3. Kachurin GA, NV Nidaev, AV Khodyachikh and LA Kovaleva : *Sov Phys Semicond* 10, p 1128, 1976.
4. Anderson CL, GK Celler and GA Rozgonyi *et al.* : *Laser and Electron Beam Processing of Electronic Materials. Electrochemical Society, New Jersey, 1979.*
5. White CW and PS Peercy *et al.* : *Laser and Electron Beam Processing of Electronic Materials. Academic Press, New York, 1980.*
6. Sealy BJ, SS Kular, KG Stephen, R Croft and A Palmer : *Electron Lett*, 14, p 721, 1978.
7. Tandon JL, MA Nicolet, WF Tseng, FH Eisen, SU Campisano, G Foti and E Rimini : *Appl Phys Lett*, 39, p 597, 1979.
8. Nojima S : *J Appl Phys*, 53, p 5028, 1982.
9. Oraby AH, Y Yuba, M Takai, K Gamo and S Namba : *Jpn J Appl Phys*, 23, p 326, 1984.
10. FH Eisen : *In Ref 5*, p 309.
11. Gamo K, Y Yuba, AH Oraby, K Murakami, S Namba and Y Kawasaki : *In Ref 5*, p 322.
12. Baeri P, SU Campisano, G Foti and E Rimini : *J Appl Phys*, 50:788, 1997.
13. Surko CM, AL Simions, DH Auston, JA Golovchenko, RE Slusher and TNC Venkatesan : *Appl Phys Lett*, 34:635, 1979.
14. Wang ZL and FW Saris : *Phys Lett*, 83A:367, 1981.
15. Wood RF and GE Giles : *Phys Rev*, B23:2923, 1981.
16. Moody JE, RH Hendel : *J Appl Phys*, 53:4364, 1982.
17. Vitali G : *Jpn J Appl Phys*, 31:2049, 1992.
18. Van Vechten JA : *In Ref 5*, p 53.
19. Muller H, FH Eisen and JW Mayer : *J Electrochem Soc*, 122:651, 1975.
20. Oraby AH, K Murakami, Y Yuba, K Gamo, S Namba and Y Masuda : *Appl Phys Lett*, 38:562, 1981.
21. Amona J, PA Pianetta and CA Stolte : *Appl Phys Lett*, 36:597, 1980.
22. Tandon JL, MA Nicolet and FH Eisen : *Appl Phys Lett*, 34:165, 1979.
23. Williams JS : *In Ref 4*, p 249.
24. Gibbons JF, WS Johnson and SW Myroie : *Projected range Statistics, 2nd ed, Johnson Wiley and Sons Inc, 1975.*
25. Surridge RK and BJ Sealy : *J Phys D:Appl Phys*, 10:911, 1977.
26. Lindow A, JF Gibbons, VR Deline and CA Evans : *Appl Phys Lett*, 32:15, 1978.
27. Nudd GR and PA Nygaard : *Elect Lett*, 14:85, 1978.
28. Sealy BJ, MH Badawi, SS Kular and KG Stephens : *In Ref 4*, p 610.
29. Gamo K, F Katano, Y Yuba, K Murakami and S Namba : *Laser-Solid Interactions and Laser Processing (1978), AIP Conf Proc No 50, edited by SD Ferris, HJ Leamy and JM Poate. Amer Inst Phys, New York, p 591, 1979.*

Correspondence:  
A. H. ORABY  
Department of Physics,  
Faculty of Science,  
Mansoura University,  
Mansoura, EGYPT.